

IM.1 Introductory Material, PDE + BCs + IC

Laplacian dominates engineering PDE descriptions

due to "*physics model*" of *continuum mechanics*

yields EBV problem statements, hence BCs, IC

Example: unsteady heat conduction with source

$$\text{DE: } \frac{\partial T}{\partial t} = \kappa \nabla^2 T + s$$

$$\text{steady, 1D form (ODE)} \quad \frac{d^2 T}{dx^2} + \frac{s}{\kappa} = 0 + \text{BCs}$$

Determine analytical solutions for:

$$\text{BCs: } T(x_L) = T_L, T(x_R) = T_R$$

$$\text{sources: } s/\kappa \Rightarrow C, \sum_i C_i x^i, \text{ and } \sum_n C_n \sin\left(\frac{n\pi x}{l}\right)$$

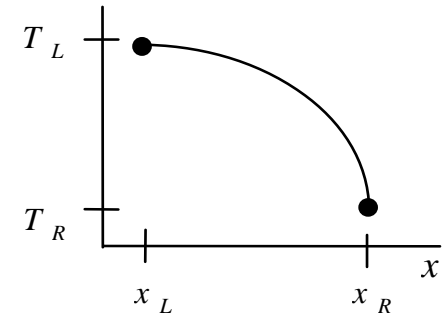
$$\text{domain: } l \equiv x_R - x_L$$

IM.2 Steady 1-D Conduction, Solutions

Constant source, integrate ODE twice

$$T(x) = \frac{1}{2} Cx^2 + ax + b$$

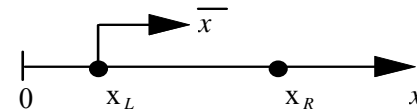
constants of integration are a and b
determine via BCs



Polynomial series source, integrate directly

$$T(x) = \frac{1}{(i+1)(i+2)} \sum_i C_i x^{i+2} + ax + b$$

translate origin, use BCs in \bar{x} system



Fourier series source, use \bar{x} system

trial space:

$$T(x) \equiv \sum_m A_m \sin\left(\frac{m\pi\bar{x}}{l}\right) + a\bar{x} + b$$

substitute into ODE

$$\Rightarrow m, n, A_m, C_n + \text{BCs } (a, b)$$

IM.3 Unsteady 1-D Conduction Solution

Conclusion: 1-D linear *laplacian* ODE is easy to solve

Greater than 1-D employs *separation of variables*

find *trial space*

determine expansion coefficients using BCs and IC

Unsteady 1-D heat conduction, $s = 0$

PDE:
$$\frac{1}{\kappa} \frac{\partial T}{\partial t} = \frac{\partial^2 T}{\partial x^2}$$

BCs:
$$T(x=0, t) = 0 = T(x=l, t)$$

IC:
$$T(x, t=0) = f(x)$$

Trial space form via separation of variables (SOV)

$$T(x, t) = \sum_{\alpha} \Psi_{\alpha}(x, t) Q_{\alpha} \equiv F(x) G(t)$$

PDE becomes:
$$\frac{1}{\kappa G} \frac{dG}{dt} = \frac{1}{F} \frac{d^2 F}{dx^2} \equiv \pm \beta^2$$

IM.4 Unsteady 1-D Conduction Solution

Separation of variables: PDE \Rightarrow 2 ODEs

functional *independence* requires the constant β^2

ODEs now have source term functional dependence!

$$\text{ODE}(x): \frac{d^2 F}{dx^2} \mp \beta^2 F = 0$$

$$\text{ODE}(t): \frac{dG}{dt} = \pm \beta^2 \kappa G$$

Must choose sign for β^2 to maintain solution boundedness

$$\text{solutions: } F(x) = A \sin(\beta x) + B \cos(\beta x)$$

$$G(t) = C \exp(-\kappa \beta^2 t)$$

IM.5 Separation of Variables, BCs

Combining F and G produces trial space members

$$\Psi_{\alpha}(x, t)Q_{\alpha} = \left[A_{\alpha} \sin(\beta x) + B_{\alpha} \cos(\beta x) \right] \exp(-\kappa\beta^2 t)$$

Contains 3 arbitrary coefficients: A_{α} , B_{α} and β^2

BCs:

$$A_{\alpha} \sin(0) + B_{\alpha} \cos(0) = 0$$

$$A_{\alpha} \sin(\beta l) + B_{\alpha} \cos(\beta l) = 0$$

non-trivial solution:

$$B_{\alpha} = 0 \text{ and } \beta \Rightarrow \beta_n = n\pi/l, \quad n = 0, 1, 2, \dots$$

Switching index label, solution *trial space* member is

$$\Psi_n(x, t) = \sin\left(\frac{n\pi x}{l}\right) \exp(-\kappa\beta_n^2 t)$$

IM.6 Separation of Variables, GWS

Trial space basis obtained via SOV process

series solution must satisfy IC, hence at $t = 0$

IC:
$$T(x, t = 0) \equiv f(x) = \sum_n^{\infty} Q_n \sin\left(\frac{n\pi x}{l}\right) \exp(0)$$

one equation for *infinite* number of unknowns Q_n

Matrix solution via *optimal* Galerkin weak statement

$$\text{GWS} \equiv \int_{\Omega} \sin\left(\frac{m\pi x}{l}\right) f(x) dx - \sum_n^{\infty} \sin\left(\frac{m\pi x}{l}\right) Q_n \sin\left(\frac{n\pi x}{l}\right) dx \equiv 0$$

for $m = 1, 2, \dots, M, \dots$

Since $\sin(n, m)$ are *orthogonal*, matrix is diagonal

$$\Rightarrow Q_n \equiv \frac{2}{l} \int_0^l \sin\left(\frac{n\pi x}{l}\right) f(x) dx$$

IM.7 Theoretical Foundations, GWS

Analytical SOV solution examples, ODEs, PDEs

- *identify the trial space for the approximation,*
- *form the expansion on Q_∞*
- *determine optimal coefficients by evaluating the GWS*
- *however, SOV infinite series PDE solution intractable !*

Resolution, truncate series for an *approximation*

$$T(x, t) \cong T^N(x, t) \equiv \sum_n^N Q_n \sin\left(\frac{n\pi x}{l}\right) \exp\left(-\kappa\beta_n^2 t\right)$$

$\sin(n)$ belongs to a *complete* set

select N for verifiable *accuracy*

Analagously, GWS ^{h} discrete implementation via FE on Ω^h

asymptotic error estimate

mesh *refinement* $\Leftrightarrow N$ increasing

IM.8 Steady 3-D Conduction SOV Solution

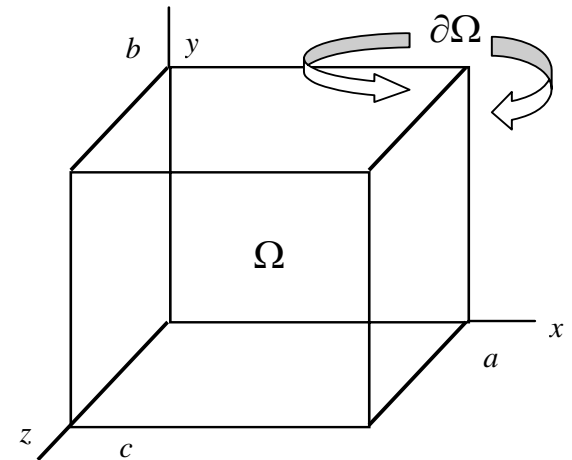
Steady conduction on 3-D cube domain

$$\text{PDE : } \nabla^2 T = \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} = 0$$

$$\text{BCs : } T(0, y, z) = 0 = T(a, y, z)$$

$$T(x, 0, z) = 0 = T(x, b, z)$$

$$T(x, y, 0) = 0, \text{ and } T(x, y, c) = f(x, y)$$



Trial space determination

$$\text{form for SOV: } T(x, y, z) = X(x) Y(y) Z(z)$$

$$\text{homogeneous BCs: } \Psi_{n,m}(\mathbf{x}) \equiv \sin\left(\frac{n\pi x}{a}\right) \sin\left(\frac{m\pi y}{b}\right) \sinh(\pi\beta_{n,m} z)$$

GWs on coefficient matching $f \Rightarrow$ diagonal matrix

$$Q_{n,m} \equiv \frac{4}{ab} \int_0^a \int_0^b f(x, y) \sin\left(\frac{n\pi x}{a}\right) \sin\left(\frac{m\pi y}{b}\right) dx dy$$

approximation:

$$T^N(\mathbf{x}) \equiv \sum_{n,m}^{N=M} \Psi_{n,m}(\mathbf{x}) Q_{n,m}$$

IM.9 Summary, Analytical PDE Methods

Restrictions for analytic SOV “to work”

PDE linear

BCs separable on $\partial\Omega$

quasi-linear data

Attribute assimilations with weak statement

solution trial space

orthogonality

completeness

predictable accuracy enhancement

Finite element implementation of GWS handles

PDE non-linear systems

BCs on arbitrary $\partial\Omega$

arbitrary data, first derivatives

IM.10 Legacy FD Connections To WS

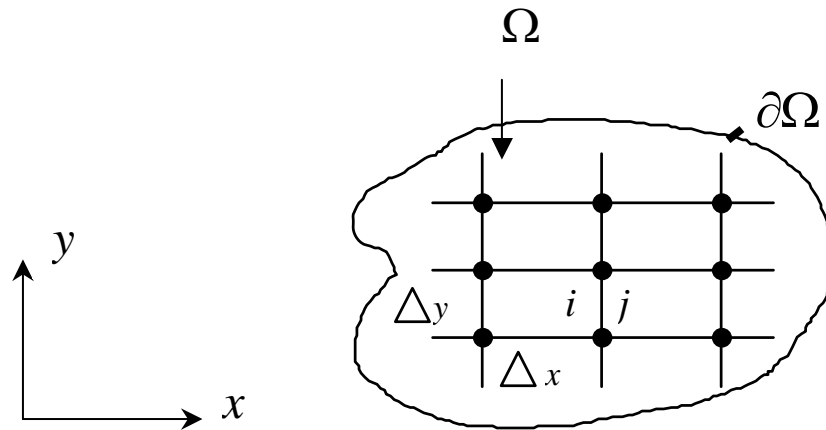
Finite difference methodology

Taylor series for PDE operators
essentially a 1-D process

Example problem:

$$L(\phi) = \nabla^2 \phi + s = 0$$

$$\ell(\phi) \Rightarrow \phi = \text{constant}$$



For cartesian interior meshing with measures Δx , Δy

$$\phi(x + \Delta x, y) = \phi(x, y) + \Delta x \frac{\partial \phi}{\partial x} \Big|_{x,y} + \frac{1}{2} \Delta x^2 \frac{\partial^2 \phi}{\partial x^2} \Big|_{x,y} + O(\Delta x^3)$$

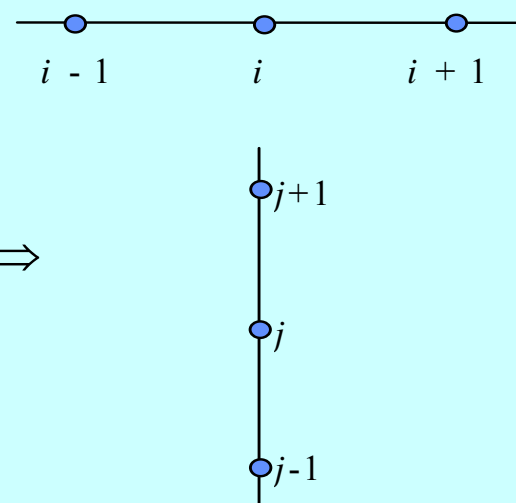
$$\Rightarrow \frac{\partial \phi}{\partial x} \Big|_{x,y} \equiv \frac{\partial \phi}{\partial x} \Big|_{i,j} = \frac{\phi_{i+1,j} - \phi_{i,j}}{\Delta x} + O(\Delta x)$$

IM.11 FD Order-Of-Accuracy Control

Backwards TS yields $O(\Delta x, \Delta y)$ constructions

$$\begin{aligned}\phi(x - \Delta x, y) &= \phi(x, y) - \Delta x \left. \frac{\partial \phi}{\partial x} \right|_{x,y} + \Delta x^2 \left. \frac{\partial^2 \phi}{\partial x^2} \right|_{x,y} + \\ &\Rightarrow \left. \frac{\partial \phi}{\partial x} \right|_{i,j} = \frac{\phi_{i,j} - \phi_{i-1,j}}{\Delta x} + O(\Delta x)\end{aligned}$$

Subtracting TSs $\Rightarrow O(\Delta x^2, \Delta y^2)$ “molecules”

$$\left. \begin{aligned}\left. \frac{\partial \phi}{\partial x} \right|_{i,j} &= \frac{\phi_{i+1,j} - \phi_{i-1,j}}{2 \Delta x} + O(\Delta x^2) \\ \left. \frac{\partial \phi}{\partial y} \right|_{i,j} &= \frac{\phi_{i,j+1} - \phi_{i,j-1}}{2 \Delta y} + O(\Delta y^2)\end{aligned} \right\} \Rightarrow$$


IM.12 FD Stencils, Laplacian

FD operators written as *stencils*

$$\left. \frac{\partial \phi}{\partial x} \right|_{i,j} = \frac{1}{2h} \begin{pmatrix} -1 & 0 & 1 \end{pmatrix} + O(h^2), \quad \left. \frac{\partial \phi}{\partial y} \right|_{i,j} = \frac{1}{2k} \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix} + O(k^2)$$

using h, k , to replace $\Delta x, \Delta y$

Via same processes 2nd derivative FD stencils

$$\left. \frac{\partial^2 \phi}{\partial x^2} \right|_{i,j} = \frac{1}{h^2} \begin{pmatrix} 1 & -2 & 1 \end{pmatrix} + O(h^2), \quad \left. \frac{\partial^2 \phi}{\partial y^2} \right|_{i,j} = \frac{1}{k^2} \begin{pmatrix} 1 \\ -2 \\ 1 \end{pmatrix} + O(k^2)$$

Hence, 2-D laplacian FD stencil (“symbol”) for $h = k$

$$\left. \nabla^2 \right|_{i,j} = \frac{1}{h^2} \begin{bmatrix} & 1 & \\ 1 & -4 & 1 \\ & 1 & \end{bmatrix} + O(h^2)$$

3-D laplacian FD symbol is the obvious extension

IM.13 FD Symbol \Leftrightarrow GWS Matrix Statement

FD symbol for laplacian PDE on uniform mesh

$$\left(\nabla^2 \phi + s \right)_{\text{FD}} \Rightarrow \frac{1}{h^2} \begin{bmatrix} 1 & & \\ & -4 & \\ & & 1 \end{bmatrix} \Phi_{i,j} + S_{i,j} + \text{TE} = 0$$

a “geometric” picture of matrix statement

Symbol correspondence to GWS statement

$$\frac{1}{h^2} \begin{bmatrix} 1 & & \\ & -4 & \\ & & 1 \end{bmatrix} \Phi_{i,j} + S_{i,j} + \text{TE} \Rightarrow [\text{Matrix}] \{ \Phi \} - \{ \mathbf{b}(s) \} + \{ \text{error} \} = \{ 0 \}$$

GWS^h generates matrix statement via FE *assembly*

FD stencil \Leftrightarrow FE [Matrix] connections will become exposed

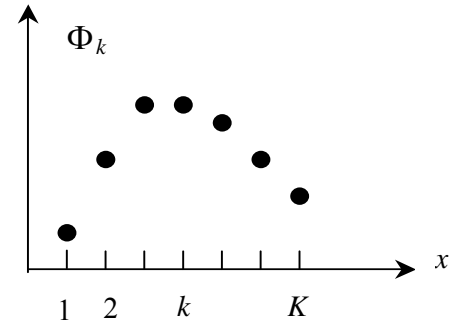
IM.14 FD Approximate Solution

FD process yields a set of numbers $\Phi_k, 1 \leq k \leq \text{nnode}$

FD numbers \Leftrightarrow approximate solution
employs interpolation concepts

Interpolation polynomial

$$f(x) = a_0 + a_1x + a_2x^2 + \dots = \sum_{\alpha=0}^N a_{\alpha} x^{\alpha}$$



$N+1$ coefficients a_{α} produce N^{th} -degree polynomial

Using nodal data, an FD "solution" could be

$$f(x_k) = \Phi_k = a_0 + a_1x_k + a_2x_k^2 + \dots + a_Nx_k^N \quad \text{for } 1 \leq k \leq K$$

\Rightarrow [Matrix] statement for determining the a_{α}

$$\begin{bmatrix} 1 & x_1 & x_1^2 & x_1^3 & \dots & x_1^N \\ 1 & x_2 & \dots & \dots & \dots & \dots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \end{bmatrix} \begin{Bmatrix} a_1 \\ a_2 \\ \vdots \\ a_N \end{Bmatrix} = \begin{Bmatrix} \Phi_1 \\ \Phi_2 \\ \vdots \\ \Phi_N \end{Bmatrix}$$

IM.15 Polynomial FD “Solution”

Evaluating polynomial at K FD nodes x_k produces

$$[\text{Matrix } (x_k^\alpha)] \{ a_k \} = \{ \Phi_k \}$$

solvability requires $K = N + 1$

involves α -exponentiation on x_k

Solution of [Matrix] statement, e.g., via Cramer’s rule

$$\text{FD solution: } \Phi|_{\text{FD}} \Rightarrow f(x) = \sum_{\substack{\alpha=0, k-1 \\ k=1, K}} a_\alpha (x_k, \Phi_k) x^\alpha$$

detraction is oscillations between nodes for $K \geq 6$

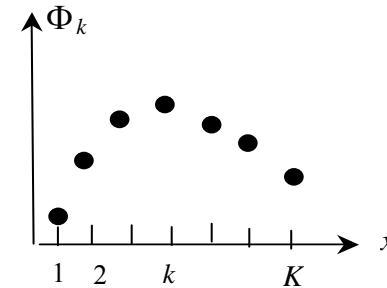
resolution is *piecewise continuous* Lagrange polynomials

IM.16 Lagrange Polynomial Interpolation

Functionally convenient interpolation polynomial

Lagrange:
$$f(x) \equiv \sum_{k=1}^K L_K(x, x_k) \Phi_k$$

rather than:
$$f(x) = \sum_{\alpha} a_{\alpha}(x_k, \Phi_k) x^{\alpha}$$



Lagrange polynomial, for K data points

$$f(x_j) \equiv \Phi_j = \sum_{k=1}^K L_K(x_j, x_k) \Phi_k = \delta_{jk} \Phi_k$$

for the $K-1$ zeros of δ_{jk}

$$L_K(x, x_k) = C_k (x - x_1) (x - x_2) \dots (x - x_{k-1}) (x - x_{k+1}) \dots (x - x_K)$$

for δ_{jk} at unity:

$$C_k = 1 / (x_k - x_1) (x_k - x_2) \dots (x_k - x_{k-1}) (x_k - x_{k+1}) \dots (x_k - x_K)$$

IM.17 Lagrange Interpolation Polynomials

Lagrange polynomial of degree $K-1$

$$L_K(x, x_k) = \frac{(x - x_1)(x - x_2) \dots (x - x_{k-1})(x - x_{k+1}) \dots (x - x_K)}{(x_k - x_1)(x_k - x_2) \dots (x_k - x_{k-1})(x_k - x_{k+1}) \dots (x_k - x_K)}$$

polynomial interpolation for FD data Φ_k

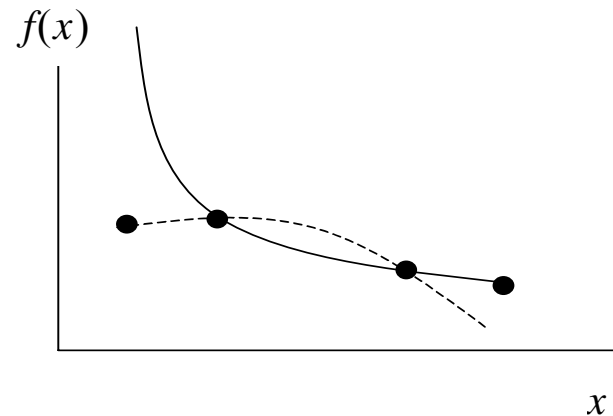
$$\phi(x)_{\text{FD}} = \sum_{k=1}^K L_K(x, x_k) \Phi_k(\text{FD})$$

provides a conceptual connection to GWS starting point

$$q^N(x) = \sum_{\alpha=1}^N \Psi_{\alpha}(x) Q_{\alpha}$$

Piecewise continuity issues

depends on K
not unique for $K > 1$



IM.18 GWS Summary, FE Trial Space Basis

FE global solution approximation, \cup denotes “union”

$$q^N \equiv q^h = \cup_e q_e(x), \text{ and } q_e(x) \equiv \{N_k(\eta)\}^T \{Q\}_e$$

key is FE trial space basis $\{N_k(\eta)\}$

construction via Lagrange polynomials
extends directly to

$$\begin{aligned} k &> 1 \\ n &> 1 \end{aligned}$$

involves local (intrinsic) coordinate system η

coordinate transformations $x = x(\eta)$

Solving GWS determines expansion coefficient set

$$[\text{Matrix}] \text{ solution} \Rightarrow \{Q\} = \cup_e \{Q\}_e$$