

FVNS.1 The Class of FVS^h-Implemented PPNS Algorithms

Commercial CFD codes for INS employ FVS^h PPNS approximations

e.g., Fluent, CFX, Star-CD, Flow 3D

Theory requirement

$$DM^h: \quad \nabla \cdot \mathbf{u} = 0 \Rightarrow \|\nabla^h \cdot \mathbf{u}^h\| \leq \varepsilon > 0 \text{ iteratively}$$

“famous” named algorithms with origination

MAC, SMAC	– Los Alamos Nat. Lab
SIMPLE,- ER, -EC, -EST	– Imperial College, UK
PISO	– Imperial College, UK

each is an iterative strategy for approximating the theory requirement

1. measure DM^h error $\nabla^h \cdot \mathbf{u}^h$ via a potential function ϕ^h
2. employ ϕ^h to moderate $D\mathbf{P}^h$ error via
 - velocity corrections
 - pressure corrections
3. iterate $D\mathbf{P}^h$ with DM^h by relaxing corrections (stability)
4. stop when $FVS^h(\mathcal{L}(q^h)) \Rightarrow |\{\text{RES}\}_{\max}| < \delta$

FVNS.2 Generic FVS^h Algorithm for PPNS Form for INS

Reynolds transport theorem generates typical $\mathcal{L}(q)$, hence

$$D(\cdot) : \mathcal{L}(q) = \frac{\partial q}{\partial t} + \frac{\partial}{\partial x_j} (f_j - f_j^r) - s = 0, \text{ on } \Omega \subset \Re^n$$

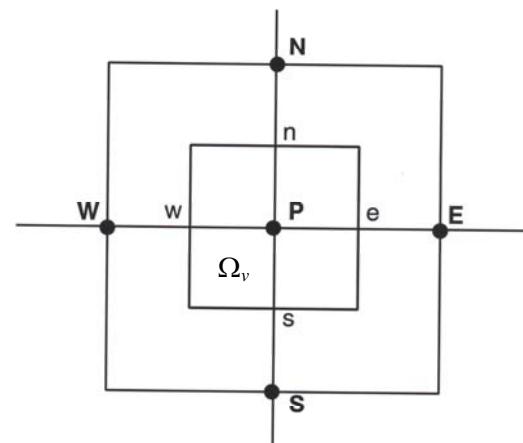
$$\begin{aligned} \text{FVS}^h &\equiv \int_{\Omega} \mathcal{L}(q^h) d\tau \equiv \sum_{\Omega^h} \int_{\Omega_v} \left(\frac{\partial q}{\partial t} - \frac{\partial}{\partial x_j} (f_j - f_j^v) - s \right) d\tau \equiv 0 \\ &= \sum_{\Omega^h} \left[\int_{\Omega_v} \left(\frac{\partial q^h}{\partial t} - s \right) d\tau - \oint_{\partial\Omega_v} (f_j - f_j^v)^h \cdot \hat{\mathbf{n}}_j d\sigma \right] = 0 \end{aligned}$$

Use of divergence theorem necessitates “staggered” meshing

generic FV notation :

$$\int_{\Omega_v} (\cdot) d\tau \Rightarrow (\cdot)_P \cdot h^n$$

$$\oint_{\partial\Omega_v} (\cdot) \cdot \hat{\mathbf{n}} d\sigma \Rightarrow \sum_{\alpha=e,w,n,s} \Delta_\alpha (\cdot)^\pm h^{n-1}$$



FVNS.3 FVS^h CFD Algorithm Notation, Fluxes

SIMPLE – type FV notation for $\mathcal{L}(q^h)$, $n = 2$: (Patankar, 1980, Ch.5)

$$\begin{aligned} \text{FVS}^h(\mathcal{L}(q^h)) \Rightarrow a_p Q_p &= a_e Q_e - a_w Q_w + a_n Q_n - a_s Q_s + S_p \\ &= \sum_{B=1}^4 a_B Q_B + S_p \end{aligned}$$

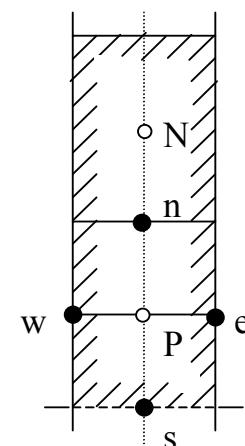
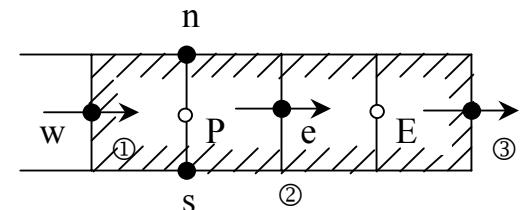
FVS^h for DP^h mixes capital and lower case node evaluations

for DP_x : $\hat{\mathbf{i}}$ – flux for P evaluated at edge e, $n = 2$

$$\text{FVS}^h(\mathcal{L}(u^h)) \Rightarrow a_e u_e = \sum_{b=1}^3 a_b u_b + (p_p - p_e) \ell_e + s_p$$

for DP_y : $\hat{\mathbf{j}}$ – flux for P evaluated at edge n, $n = 2$

$$\text{FVS}^h(\mathcal{L}(v^h)) \Rightarrow a_n v_n = \sum_{b=1}^3 a_b v_b + (p_p - p_n) \ell_n + s_p$$



FVNS.4 FVS^h Pressure-Velocity Corrections

FVS^h solution for $(u, v, w)^h$ does not satisfy $DM^h \Rightarrow \|\nabla^h \cdot \mathbf{u}^h\| < \varepsilon$

predictor-corrector strategy:

$$\mathbf{u}^{n+1} = \mathbf{u}^h + \mathbf{u}'$$

$$p^{n+1} = p^h + p'$$

DP^h SIMPLE-type correction procedure

solve : $\nabla^2 p' = \nabla \cdot \mathbf{u}^h$, for BCs $\hat{\mathbf{n}} \cdot \nabla \phi = 0$

update pressure : $p^{n+1} \equiv p^h + p'$

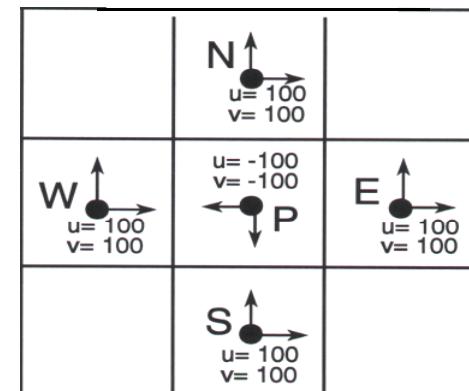
correct velocity : $\mathbf{u}^{n+1} = \mathbf{u}^h + \nabla p'$

under-relax \mathbf{u}', p' , resolve FVS^h($\mathcal{L}(\mathbf{u}^h)$) until \mathbf{u}', p' small enough

Central differences can produce spurious DM^h

example: $\nabla^h \cdot \mathbf{u}_p = 0$ for $2\Delta x$ error mode

leads to upwinding, QUICK differencing
dissipate dispersion error



FVNS.5 FVS^h SIMPLE-Type PPNS Implementations

Recall PPNS analytical theory ingredients

$$DM^h : \mathbf{u}^{n+1} - \mathbf{u}^h = -\nabla \phi$$

$$\text{divergence: } \mathcal{L}(\phi) = -\nabla^2 \phi + \nabla \cdot \mathbf{u}^h = 0$$

$$\text{BCs: } \mathcal{L}(\phi) = \hat{\mathbf{n}} \cdot \nabla \phi = 0, \quad \phi_{\text{out}} = 0$$

$$\text{DP}^h \text{ pressure: } \text{SUM}\Phi_{n+1}^{p+1} = \Phi_n + (\theta \Delta t)^{-1} \sum_{\alpha=0}^p \phi_{n+1}^{\alpha+1}$$

$$\text{genuine pressure: } \mathcal{L}(p) = -\text{Eu} \nabla^2 p - s(\nabla \mathbf{u}, \Theta) = 0$$

$$\ell(p) = \hat{\mathbf{n}} \cdot \nabla p + f(\text{Re}, \nabla^2 \mathbf{u} \cdot \hat{\mathbf{n}})$$

$$\text{accuracy assessment: } |\phi^{p+1}|_E \ll \varepsilon, \quad |\text{SUM}\Phi|_E < \varepsilon \quad \text{and} \quad |p^h|_E > \varepsilon$$

SIMPLE implementations readily interpreted within the theory

genuine pressure *never* determined from $\mathcal{L}(p)$

$\Rightarrow \text{SUM}\Phi$ is interpreted as pressure

correction pressure is identical to ϕ^p

velocity correction utilizes DM^h definition

$\Rightarrow \mathbf{u}' \cong \nabla^h \phi^h$

FVNS.6 FVS^h SIMPLE-Type PPNS Implementations, Con'd

SIMPLE (Spalding, et al, 1968)

1. guess p_{n+1}^0
2. solve $\text{d) } \Rightarrow \mathbf{u}^h$
3. solve $\text{a) } \Rightarrow p'$
4. relaxing p' , update $\text{b) } \Rightarrow p_{n+1}^1$
5. evaluate $\text{c) } \Rightarrow \mathbf{u}'$
6. repeat 1–5 until $\text{d) } |\text{RES}|_{\max} < \delta$
7. advance index n

PPNS theory ingredients

- a) $\mathcal{L}(\phi) = -\nabla^2 \phi + \nabla \cdot \mathbf{u}^h = 0$
- b) $\text{SUM } \Phi_{n+1}^p = \Phi_n + \Delta t^{-1} \sum_a^p \phi^a$
- c) $-\nabla \phi^p = \mathbf{u}^{n+1} - \mathbf{u}^h$
- d) $\text{DP}^h \Rightarrow \text{FVS}^h(\mathcal{L}(\mathbf{u}^h)) = \{\text{RES}\}_{n+1}$

SIMPLER (Patankar, 1981)

1. guess \mathbf{u}_{n+1}^0
2. for $p_{n+1}' \equiv 0$, solve $\text{d) } \Rightarrow \mathbf{u}_{n+1}^1$
3. solve $\text{a) } \Rightarrow p'$
4. update $\text{b) } \Rightarrow p_{n+1}^1$
5. solve $\text{d) } \Rightarrow u_{n+1}^2$
6. solve $\text{a) } \Rightarrow p'$
7. evaluate $\text{c) } \Rightarrow \mathbf{u}''$
8. repeat 1–7 until $\text{d) } |\text{RES}|_{\max} < \delta$
9. advance index n

Comments on performance

0. SIMPLE \Rightarrow Semi-Implicit Method for Pressure-Linked Equations
1. all employ linear stationary iteration e.g., Gauss-Seidel
2. SIMPLE slow to converge for p hence SIMPLER
3. SIMPLEC (Rathby, et. al., 1984) SIMPLE iteration with refined velocity update $\text{c) } \Rightarrow \mathbf{u}''$

FVNS.7 FVS^h PISO PPNS Implementation, Genuine Pressure

PISO \Rightarrow Pressure-Implicit with Splitting of Operators (Issa, 1984)

A non-iterative predictor-corrector

1. using p_n solve d) $\Rightarrow \mathbf{u}_{n+1}^1$
2. solve a) $\Rightarrow p'$
3. evaluate b) $\Rightarrow p_{n+1}^1$
4. evaluate c) $\Rightarrow \mathbf{u}'$
5. solve e) $\Rightarrow p_{n+1}^2$
6. solve d) $\Rightarrow \mathbf{u}_{n+1}^2$
7. advance index n

PPNS theory ingredients

- a) $\mathcal{L}(\phi) = -\nabla^2\phi + \nabla \cdot \mathbf{u}^h = 0$
- b) $\text{SUM}\Phi_{n+1}^p = \Phi_n + \Delta t^{-1} \sum_a^p \phi^a$
- c) $-\nabla\phi^p = \mathbf{u}^{n+1} - \mathbf{u}^h$
- d) $\mathbf{DP}^h \Rightarrow \text{FVS}^h(\mathcal{L}(\mathbf{u}^h)) = \{\text{RES}\}_{n+1}$
- e) $\nabla \cdot \mathbf{DP}^h \Rightarrow \text{FVS}^h(\mathcal{L}(p^h, \text{Re}))$

Comments on PISO

1. $\nabla \cdot \mathbf{DP}^h \Rightarrow \mathcal{L}(p)$ with BCs $\hat{\mathbf{n}} \cdot \nabla p = 0$
requires 3rd derivatives of \mathbf{u}^h in $\{\text{RES}\}!$
2. accepts step 6 solution as converged
 $\Rightarrow |\phi^h|_E < ?$

FVNS.8 FVS^h for PPNS System, Stability, Numerical Diffusion

FVS^h commercial CFD codes based on PPNS employ numerical diffusion

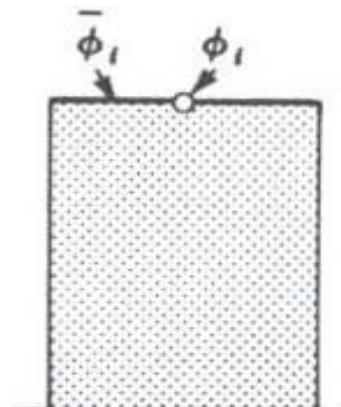
options for altering symmetric evaluation of $\int_{\partial\Omega_v} (u_j u_i) \hat{\mathbf{n}}_j d\sigma$

- direct upwinding \Rightarrow TS accuracy is $O(h)$
- hybrid upwind $\Rightarrow O(h \leftrightarrow h^2)$
- quadratic upwind $\Rightarrow O(h^3)$

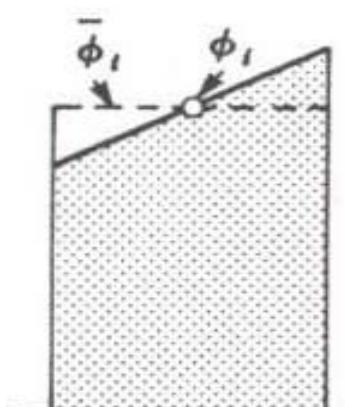
these typically employ $\theta = 1$ equivalent time – differencing

- time – accurate requires $\theta = 0.5$

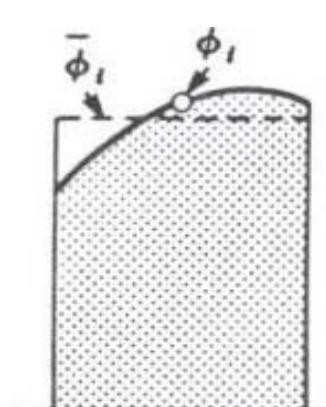
Formulation options based on Ω_v , interpolation (B. Leonard, *Adv. N.H.T.*, 1997)



constant



linear



quadratic

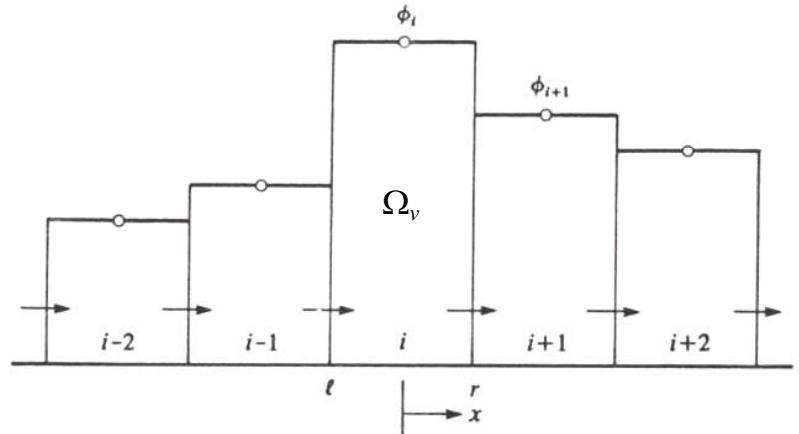
FVNS.9 Convective Flux Differencing for FVS^h PPNS

First-order upwind:

on Ω_v : $\phi = \text{constant}$

$$\oint_{\partial\Omega_v} \mathbf{u}\phi \cdot \hat{\mathbf{n}} d\sigma = (u_r \phi_r - u_\ell \phi_\ell)$$

$$\Rightarrow \phi_r (\text{1UP}) = \phi_i, u_r > 0 \\ = \phi_{i+1}, u_r < 0$$

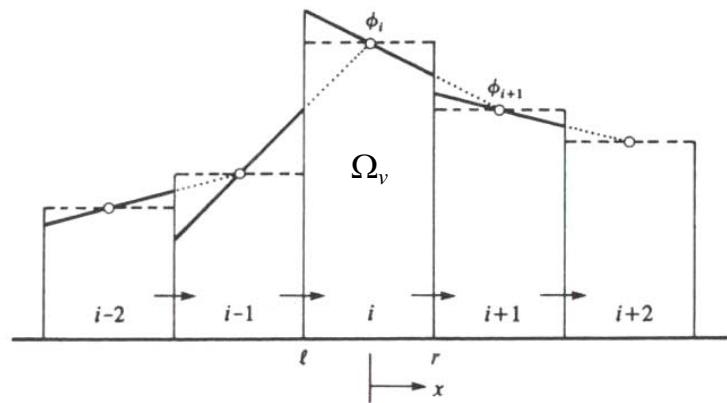


Second-order central:

on Ω_v : $\phi(x) = \phi_i + (\phi_{i+1} - \phi_i)x/h, u_r > 0$

$$\int_{\partial\Omega_r} \mathbf{u}\phi \cdot \hat{\mathbf{i}} = \phi_r (2\text{CE}) = \phi(x = h/2 - \varepsilon)$$

$$= \frac{1}{2}(\phi_{i+1} - \phi_i), u_r > 0$$



FVNS.10 Convective Flux Differencing Methods for FVS^h PPNS

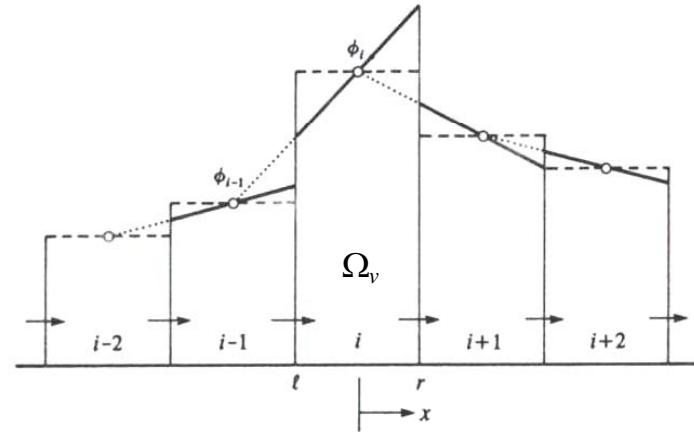
Second-order upwind:

on Ω_v : $\phi(x) = \phi_i + (\phi_i - \phi_{i-1})x/h, u_r > 0$

$$\int_{\partial\Omega_r} \mathbf{u}\phi \cdot \hat{\mathbf{i}} = \phi_r (2\text{UP}) = \phi(x = h/2 - \varepsilon)$$

$$= \frac{3}{2}\phi_i - \phi_{i-1}, \quad u_r > 0$$

$$= \frac{3}{2}\phi_{i+1} - \phi_{i+2}, \quad u_r < 0$$



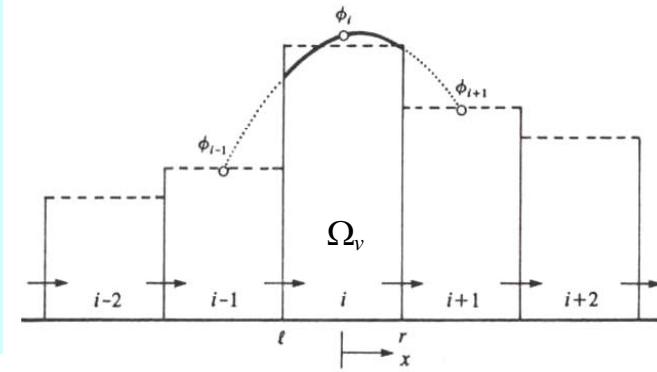
Third-order upwind (QUICK):

on Ω_v : $\phi(x) = \phi_i + \frac{1}{2}(\phi_{i+1} - \phi_{i-1})x/h + \frac{1}{2}(\phi_{i+1} - 2\phi_i + \phi_{i-1})(x/h)^2$

$$\int_{\partial\Omega_r} \mathbf{u}\phi \cdot \hat{\mathbf{i}} = \phi_r (3\text{QK}) = \phi(x = h/2 - \varepsilon)$$

$$= (6\phi_i + 3\phi_{i+1} - \phi_{i-1})/8, \quad u_r > 0$$

$$= (3\phi_i + 6\phi_{i+1} - \phi_{i+2})/8, \quad u_r < 0$$



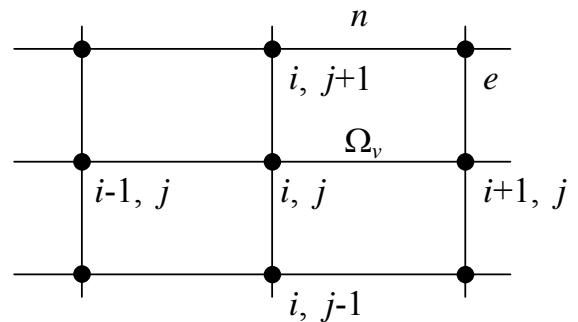
FVNS.11 QUICK Differencing for Convection, $n = 1, 2$

QUICK upwind differencing in curvature factor form, $n = 1$

$$\begin{aligned}\phi_r(3QU) &= \frac{1}{2}(\phi_{i+1} + \phi_i) - \frac{1}{8}(\phi_{i+1} - 2\phi_i + \phi_{i-1}), \quad u_r > 0 \\ &= \frac{1}{2}(\phi_{i+1} + \phi_i) - \frac{1}{8}(\phi_{i+2} - 2\phi_{i+1} + \phi_i), \quad u_r < 0 \\ &= \text{average } \phi|_{\partial\Omega_r} + c\delta^2\phi|_{\Omega_v}\end{aligned}$$

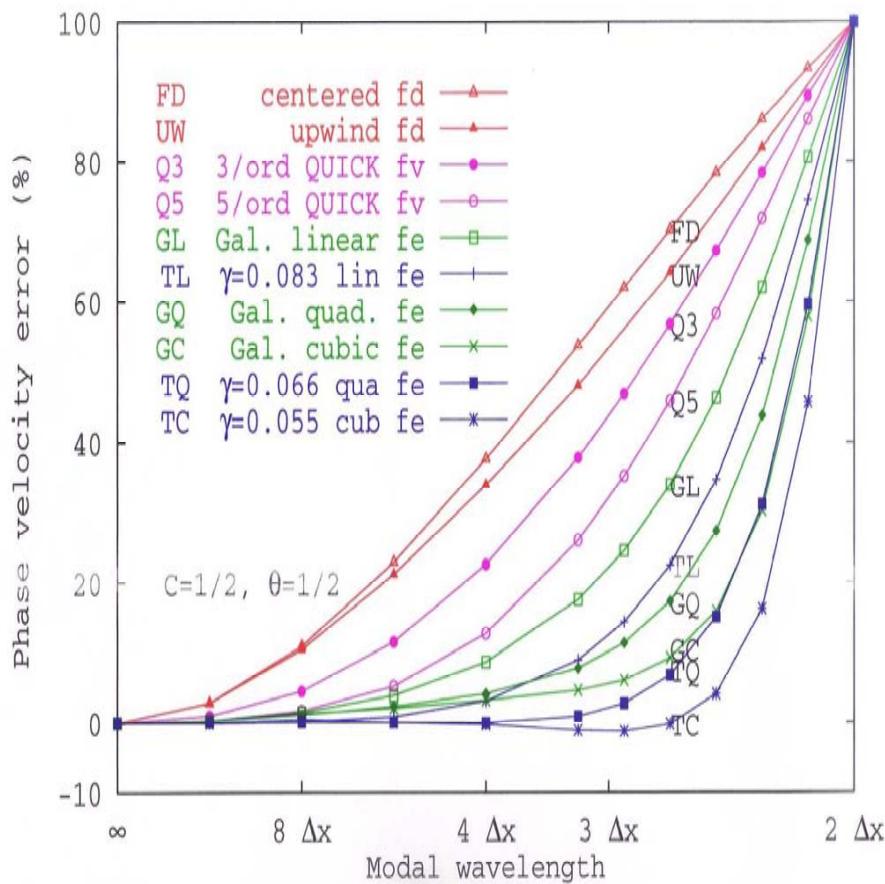
Multi-dimensional QUICK differencing for convection

$$\begin{aligned}\phi_{east}(3QU) &= \frac{1}{2}(\phi_{i+1,j} + \phi_{i,j}) \\ &\quad - \frac{1}{8}(\phi_{i+1,j} - 2\phi_{i,j} + \phi_{i-1,j}) \\ &\quad + \frac{1}{24}(\phi_{i,j+1} - 2\phi_{i,j} + \phi_{i,j-1}), \quad u_{east} > 0 \\ &= \text{average } \phi|_{\partial\Omega_v} + c_1\delta_x^2\phi|_{\Omega_v} - c_2\delta_y^2\phi|_{\Omega_v}\end{aligned}$$

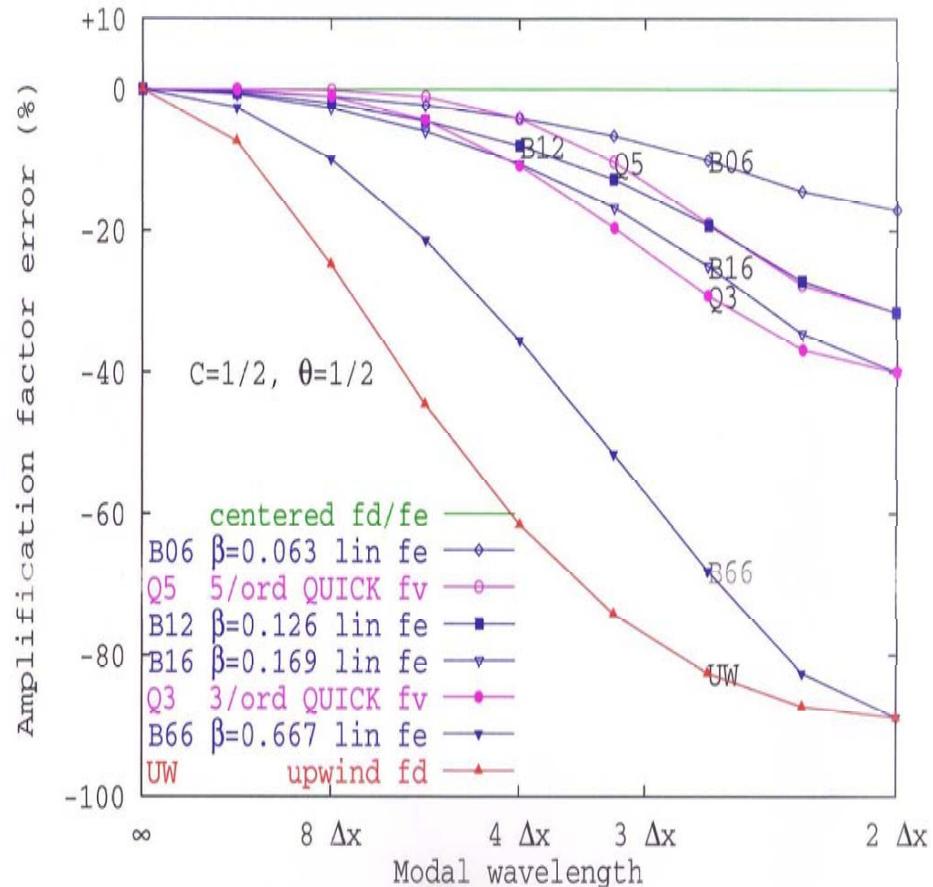


FVNS.12 GWS^h-TWS^h-FVS^h Spectral Resolution Comparisons

Phase velocity error



Artificial diffusion error



FVNS.13 Summary FVS^h - Implemental PPNS Algorithms

PPNS analytical theory ingredients are independent of implementation

$$DM^h : \mathbf{u}^{n+1} - \mathbf{u}^h = -\nabla \phi$$

$$\text{divergence} : \mathcal{L}(\phi) = -\nabla^2 \phi + \nabla \cdot \mathbf{u}^h = 0$$

$$\text{BCs} : \mathcal{L}(\phi) = \hat{\mathbf{n}} \cdot \nabla \phi = 0, \quad \phi_{\text{out}} = 0$$

$$\text{DP}^h \text{ pressure} : \text{SUM } \Phi_{n+1}^{p+1} = \Phi_n + (\theta \Delta t)^{-1} \sum_{\alpha=0}^p \phi_{n+1}^{\alpha+1}$$

$$\text{genuine pressure} : \mathcal{L}(p) = -\text{Eu} \nabla^2 p - s(\nabla \mathbf{u} \mathbf{u}, \Theta) = 0$$

$$\ell(p) = \hat{\mathbf{n}} \cdot \nabla p + f(\text{Re}, \nabla^2 \mathbf{u} \cdot \hat{\mathbf{n}})$$

$$\text{accuracy assessment} : \left| \phi^{p+1} \right|_E \ll \varepsilon, \quad \left| \text{SUM } \Phi \right|_E < \varepsilon \quad \text{and} \quad \left| p^h \right|_E > \varepsilon$$

FVS^h implementations constitute substantial simplifications

genuine pressure *never* determined from $\mathcal{L}(p) + \ell(p)$

employ stationary linear iteration predictor–corrector procedures

$\nabla^h \phi^h \Rightarrow$ velocity correction

$\phi^h \Rightarrow$ pressure correction

SUM $\Phi \Rightarrow$ interpreted as genuine pressure

convergence estimate based on $\max \left| \text{RES} \right|^p$

$\left| \phi^h \right|_E$ is never evaluated

artificial diffusion permeates implementations