

TFD.1 Fluid Dynamics, the Physics of Turbulence

Eulerian conservation principles $DM + DP$

basic scales are length and velocity $\Rightarrow Re = UL/\nu$

$Re < \delta$, $DM + DP \Rightarrow$ stable, time – independent solutions

$Re > \delta$, $DM + DP \Rightarrow$ time dependent solutions characterized by

$\left. \begin{array}{l} \text{cascades of vortices} \\ \text{cascade of length scales} \end{array} \right\} \text{"turbulence"}$

Turbulent fluid motion

“turbulent fluid motion is an irregular condition of flow in which the various measures show a random variation in time and space such that statistically distinct average values can be discerned,” (Hinze, 1975)

$\Rightarrow$ turbulence is a *continuum* phenomena
develops as an instability of laminar flow

TFD.2 Fluid Dynamics, the Physics of Turbulence

General properties of turbulence

- length scales range from domain size (L) to small (η)
- generates appearance of an “eddying” motion
- large scale (L) motion carries most of the energy
responsible for enhanced diffusion processes
- smallest scale (η) motion $\gg$ molecular mean free path
most vorticity resides in η scale motion
- turbulence energy cascades from $L \Rightarrow \eta$ scales
turbulent flows are always dissipative
- turbulent flow state depends on upstream flow history
local strain-rate tensor description not unique

TFD.3 Fluid Dynamics, the Physics of Turbulence

Kolmogorov scales in turbulent flow

all turbulent flows transfer turbulence kinetic energy k from L to smaller scales

conversion to heat via viscosity ν occurs only at the smallest scales

dissipation at the smaller scales depends only on

$\varepsilon = -dk/dt$, larger eddy energy supply rate

$\nu =$, fluid kinematic viscosity

Kolmogorov scales of turbulence

for: $\varepsilon = D(\text{length}^2 / \text{time}^3)$, $\nu = D(\text{length}^2 / \text{time})$

length: $\eta = (\nu^3 / \varepsilon)^{1/4}$

time: $\tau = (\nu^3 / \varepsilon)^{1/2}$

velocity: $v = (\nu \varepsilon)^{1/4}$

note: at STP, $\eta \approx O(10^2)$ times molecular mean free path

TFD.4 Fluid Dynamics, the Physics of Turbulence

The energy spectrum of turbulent flows

spectrum variable is wave number κ

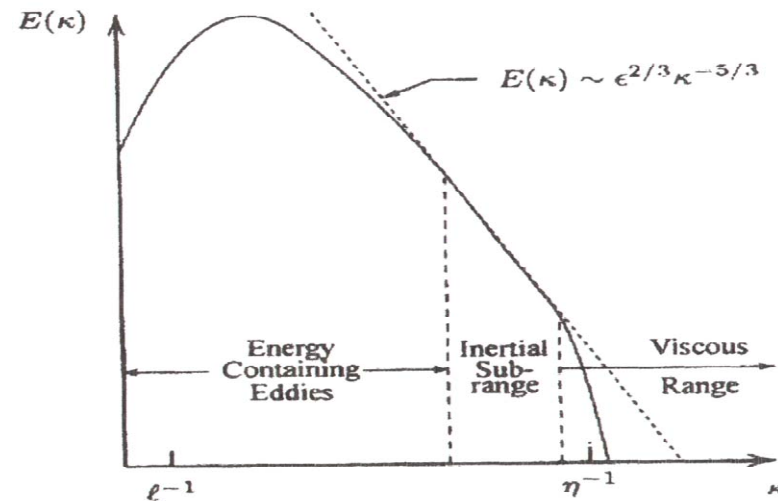
$\lambda = 2\pi/\kappa$ is wave length

$\propto$ "eddy size"

turbulence kinetic energy

$$k \equiv \int_0^\infty E(\kappa) d\kappa$$

$E(\kappa)$ is energy spectral density
(related to Fourier transform)



Dimensional analysis leads to the following observations / correlations

$$\varepsilon \sim k^{3/2} / \ell \quad \text{and} \quad \ell \gg \eta$$

$$\frac{\ell}{\eta} = \frac{\ell}{(v^3/\varepsilon)^{1/4}} \sim \frac{\ell (k^{3/2}/\ell)^{1/4}}{v^{3/4}} \sim \text{Re}_T^{3/4}, \quad \text{Re}_T \equiv \frac{k^{1/2} \ell}{v} \gg 1$$

$$\therefore E(\kappa) \Rightarrow E(\varepsilon, \kappa) = C_\kappa \varepsilon^{2/3} \kappa^{-5/3}, \quad \frac{1}{\ell} \ll \kappa \ll \frac{1}{\eta}$$

TFD.5 Fluid Dynamics, the Physics of Turbulence

Perturbation theory removes arbitrariness of dimensional analysis

inertial region: $E(\kappa) \Rightarrow E(\varepsilon, \kappa, \nu) = \varepsilon^{1/4} \nu^{5/4} f(\kappa \eta)$

outer region: $E(\kappa) = k \ell g(\kappa \ell)$ in small wave number region

matching: $\varepsilon^{1/4} \nu^{5/4} f(\kappa \eta) = k \ell g(\kappa \ell)$, $\kappa \eta \ll 1$, $\kappa \ell \gg 1$

"between viscous and energetic turbulent flow scales an overlap region must exist wherein solutions match as $\text{Re} \Rightarrow \infty$ "

enforcing smooth matching in wave number space, $d(\cdot)/d\kappa$ must be continuous

$$\Rightarrow \eta \varepsilon^{1/4} \nu^{5/4} f'(\kappa \eta) = \ell^2 k g'(\kappa \ell), \quad \kappa \eta \ll 1, \quad \kappa \ell \gg 1$$

recalling: $\eta = \nu^{3/4} \varepsilon^{-1/4}$, $k = \varepsilon^{2/3} \ell^{2/3}$

$$\Rightarrow (\kappa \eta)^{8/3} f'(\kappa \eta) = (\kappa \ell)^{8/3} g'(\kappa \ell) = \text{constant}$$

since: $\kappa \eta$ and $\kappa \ell$ are independent variables

integrating

$$\Rightarrow f(\kappa \eta) = C_\kappa \varepsilon^{2/3} \kappa^{-5/3}$$

TFD.6 Fluid Dynamics, Turbulent BL Similarity

Turbulent boundary layer near wall data exhibits self-similar log profile

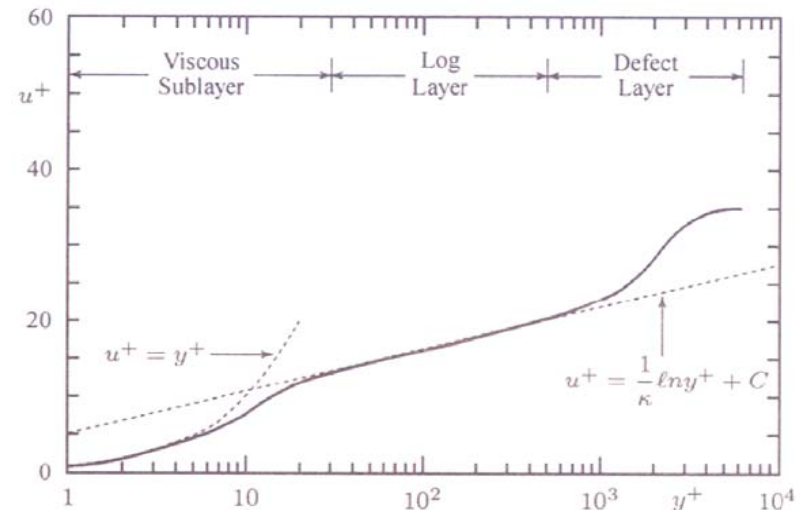
law – of – the – wall : $U/u_\tau = \kappa^{-1} \ln y^+ + C$, $y^+ = u_\tau y / \nu$

Dimensional analysis provides insight

momentum exchange to wall $\Rightarrow \tau$
 near wall $\tau(y) \approx \text{constant} \Rightarrow \tau_w$
 similar profiles for liquids and gases

$$\frac{\tau_w}{\rho} \equiv u_\tau^2 \Rightarrow D(\text{length}^2 / \text{time}^2)$$

$$\frac{\mu}{\rho} = \nu \Rightarrow D(\text{length}^2 / \text{time})$$



mean flow profile near wall, for velocity and length scales $u_\tau = \sqrt{\tau_w / \rho}$ and ν / u_τ

$$\frac{\partial U}{\partial y} = \frac{u_\tau}{y} F(u_\tau y / \nu) = \frac{u_\tau}{y} F(y^+)$$

data correlations confirm $F(u_\tau y / \nu) \Rightarrow \kappa^{-1}$ as ν influence vanishes, i.e., $y^+ \Rightarrow \infty$

integrating : $\frac{U(y)}{u_\tau} = \frac{1}{\kappa} \ln \left(\frac{u_\tau y}{\nu} \right) + C$, $C \cong 5.0$, $\kappa \cong 0.41$

TFD.7 Fluid Dynamics, Turbulent BL Characterization

Turbulent BL data confirm three distinct “flow physics” regions

viscous sub – layer : $u^+ \approx y^+$, $1 \leq y^+ < 5$

log layer : $u^+ = \kappa^{-1} \ln(y^+) + C$, $30 \leq y^+ \leq 600$

defect layer : u^+ not correlated by $\ln(y^+)$, $y^+ > 600$

Law-of-the-wall intercept depends on wall roughness

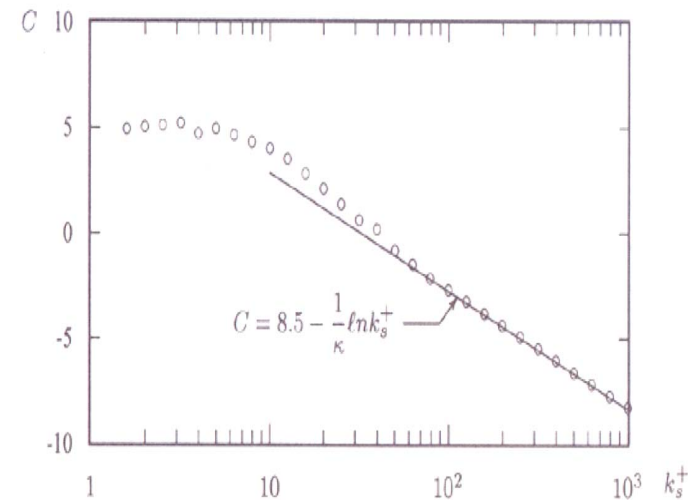
non – D roughness height : $k_s^+ \equiv \frac{u_\tau k_s}{\nu}$

Nikuradse data : $C \rightarrow 8.5 - \kappa^{-1} \ln(k_s^+)$, $k_s^+ \gg 1$

completely rough wall : $u^+ = \frac{1}{\kappa} \ln\left(\frac{y}{k_s}\right) + 8.5$

note : is independent of viscosity ν

has impact when considering LES BCs



TFD.8 Fluid Dynamics, Turbulent BL Characterization

Boundary layer defect layer lies between log layer and BL edge (δ)

$$\text{data correlation : } u^+ = \frac{1}{\kappa} \ln(y^+) + C + \frac{2\Pi}{\kappa} \sin^2\left(\frac{\pi}{2} \frac{y}{\delta}\right)$$

Cole's wake parameter Π :

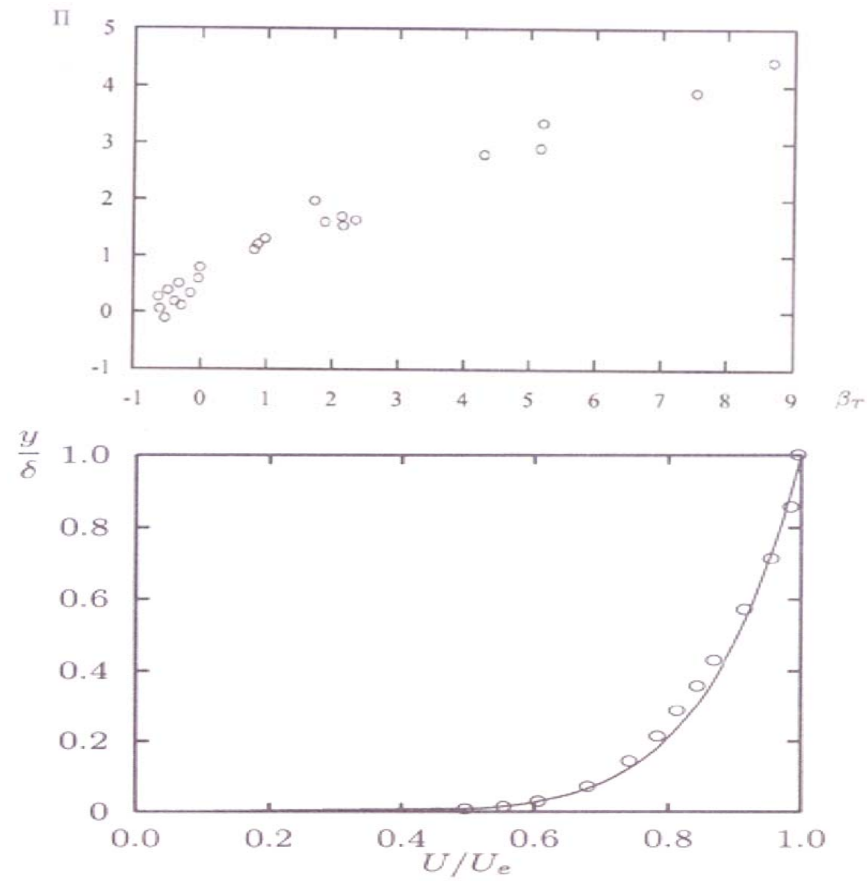
$$\beta_T \equiv \frac{\delta^* dP}{\tau_w dx}$$

$\delta^* \equiv$ BL displacement thickness

Power law approximation

turbulent BL interpolation:

$$\frac{U}{U_e} \cong \left(\frac{y}{\delta}\right)^{1/n}, \text{ for } n = 7 \rightarrow$$



TFD.9 Fluid Dynamics, Conservation Principle Manipulations

Velocity field resolved into mean and fluctuating components

resolution: $u_i(\mathbf{x}, t) \equiv U_i(\mathbf{x}, t) + u'_i(\mathbf{x}, t)$

time – averaging for stationary turbulent flow

$$\bar{u}_i(\mathbf{x}) \equiv U_i(\mathbf{x}) \equiv \lim_{T \rightarrow \infty} \frac{1}{T} \int_t^{t+T} u_i(\mathbf{x}, \tau) d\tau$$

$$\bar{u}'_i \equiv \lim_{T \rightarrow \infty} \frac{1}{T} \int_t^{t+T} (u_i(\mathbf{x}, \tau) - U_i(\mathbf{x}, \tau)) d\tau = 0$$

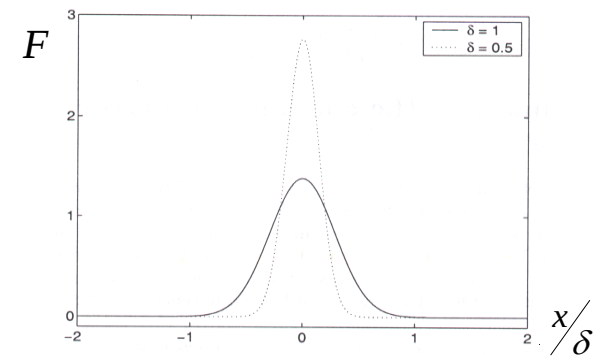
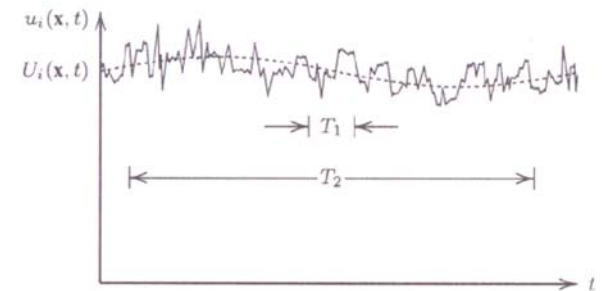
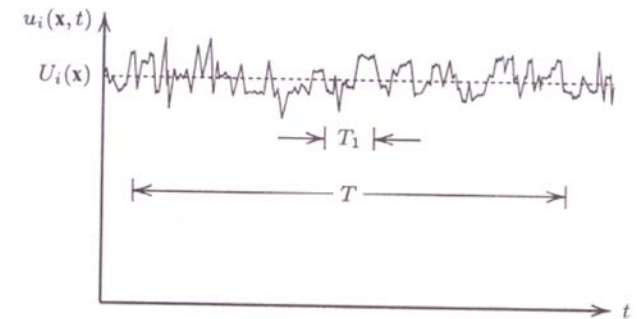
time – averaging for non – stationary turbulent flow

$$u_i(\mathbf{x}, t) = \frac{1}{T} \int_t^{t+T} u_i(\mathbf{x}, \tau) d\tau, \quad T_1 \leq T \leq T_2$$

space – filtering for non – stationary turbulence

$$\bar{f}(\mathbf{x}, t) = F * f(\mathbf{x}, t) \equiv \int_{\Omega} F(\mathbf{x} - \mathbf{z}) f(\mathbf{z}, t) d\mathbf{z}$$

$$\text{gaussian filter: } F_i(x_i) \equiv \sqrt{\frac{\gamma}{\pi \delta^2}} \exp\left(\frac{-\gamma}{\delta^2} x_i^2\right)$$



TFD.10 Fluid Dynamics, Conservation Principle Manipulations

Convection term in DP leads to Reynolds stress tensor

time – averaging : $(\overline{\mathbf{u}\mathbf{u}})_t \Rightarrow \overline{\mathbf{u}\mathbf{u}} + \overline{\mathbf{u}'\mathbf{u}'}$

space - filtering : $(\mathbf{u}\mathbf{u})_\Omega \Rightarrow \overline{\mathbf{u}\mathbf{u}} + \overline{\mathbf{u}\mathbf{u}'} + \overline{\mathbf{u}'\mathbf{u}} + \overline{\mathbf{u}'\mathbf{u}'}$

Closure requirement is to correlate fluctuating averages to state variable

MLT algebraic : $\overline{u'v'} \cong v^t \frac{\partial U}{\partial y}$, $v^t \equiv (\omega \ell)^2 \left| \frac{\partial U}{\partial y} \right|$

PNS algebraic : $\overline{u'_i u'_j} \cong \frac{2}{3} k \delta_{ij} - v^t S_{ij}$, $S_{ij} = U_{i,j} + U_{j,i}$

TKE differential : $L(k, \varepsilon) = 0$ added, $v^t = C_\mu k^2 / \varepsilon$

Turbulence intensities

flat plate BL

$$k = \frac{1}{2} \left[\overline{u'u'} + \overline{v'v'} + \overline{w'w'} \right]$$

